

Multiple hypothesis testing and testing one hypothesis multiple times: two sides of the same coin?

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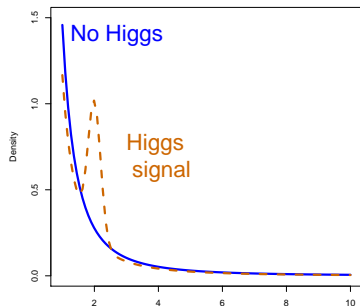
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General framework

Goal of statistical signal detection in physics

We would like to distinguish signals of new physics phenomena from the random fluctuations of the data.



- **E.g.**, Higgs boson, quark, neutrino.
- We want to detect a bump (the signal of the new particle) on top of a background flux.

How does statistics tackle this problem?

- Approach 1:
Multiple hypothesis testing \Rightarrow Bonferroni's correction.
- Approach 2:
Simulations \Rightarrow Monte Carlo, Bootstrap.
- Approach 3:
Hypothesis testing when a nuisance parameter is present only under the alternative \Rightarrow Davies (1977, 1987), Gross and Vitells (2010).



We refer to this as **Testing one hypothesis multiple times.**

Note!

In High Energy Physics a discovery is claimed at 5σ significance
 \Rightarrow in Approach 2 we need to simulate $O(10^8)$, **can we avoid that?**

Yes! Use (responsibly) Approach 1 and/or Approach 3.

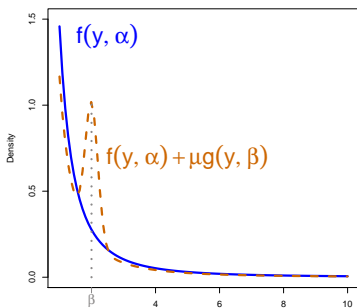
Questions I would like to address with this talk

- 1 What does it mean exactly to “test one hypothesis multiple times”, and in what sense is it equivalent to a testing problem when a nuisance parameter is present only under the alternative?
- 2 Can we tackle both nested and non-nested models with this approach?
- 3 What is the difference between testing one hypothesis multiple times and multiple hypothesis testing?
- 4 When do multiple hypothesis testing and testing one hypothesis multiple times coincide in some sense?
- 5 What else can we do, and what is the potential of working in this direction?

Outline

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A statistical framework for a physics problem



The model of interest is proportional to

$$\underbrace{f(y, \alpha)}_{\text{background}} + \underbrace{\mu}_{\text{signal strength}} \underbrace{g(y, \beta)}_{\substack{\text{signal} \\ \text{location} \\ \text{bump}}} \quad (1)$$

and we test

$$H_0 : \mu = 0 \quad \text{vs.} \quad \mu > 0. \quad (2)$$

Problems

μ is on the boundary of its parameter space + β is not defined under H_0 .

Solutions

Chernoff, 1954 + Davies, 1987, Gross and Vitells, 2010.

Theoretical solutions

Practical solution

Testing on the boundary of the parameter space

- **Model:**

$$\propto f(y, \alpha) + \mu g(y, \beta) \quad \mu \geq 0 \quad (3)$$

For now, let β be fixed, the model in (3) is identifiable.

- **Test**

$$H_0 : \mu = 0 \quad \text{versus} \quad H_1 : \mu > 0$$

- **Test statistics*:**

$$LRT = -2 \log \left[\underbrace{L(0, \hat{\alpha}_0, -)}_{\text{Likelihood under } H_0} - \underbrace{L(\hat{\mu}, \hat{\alpha}, \beta)}_{\text{Likelihood under } H_1} \right] \quad (4)$$

* for the specific case of β be fixed.

- η is on the boundary \Rightarrow WE CAN USE Chernoff, 1954 i.e.:

$$LRT \xrightarrow[n \rightarrow \infty]{d} \frac{1}{2} \chi_1^2 + \frac{1}{2} \delta(0) \quad \text{under } H_0 \quad (5)$$

Testing one hypothesis multiple times (1)

- If β fixed, under H_0 the LRT is asymptotically $\frac{1}{2}\chi_1^2 + \frac{1}{2}\delta(0)$.
- If we let β vary \Rightarrow Under H_0 , $\{LRT(\beta), \beta \in \mathbf{B}\}$ is asymptotically a $\frac{1}{2}\chi_1^2 + \frac{1}{2}\delta(0)$ random process indexed by β .
- In practice:
 - Define a grid \mathbf{B}_R of R β_r values over the energy spectrum \mathbf{B} .
 - $\forall \beta_r \in \mathbf{B}_R$ calculate $LRT(\beta_r)$.

Many “sub”-alternatives...

It is like if we had many alternative hypothesis $H_{11}, \dots, H_{1r}, \dots, H_{1R}$, one for each value $\beta_r \in \mathbf{B}_R$, and for each of them we have one value $LRT(\beta_r)$.

...but yet just one test statistic...

We finally combine the R $LRT(\beta_r)$ values in a unique test statistics

$$\max_{\beta_r \in \mathbf{B}_R} LRT(\beta_r)$$

Testing one hypothesis multiple times (2)

... and one global p-value...

The **p-value** of our test $H_0 : \eta = 0$ versus $H_a : \eta > 0$ is in the form

$$P(\sup_{\beta \in \mathbf{B}} LRT(\beta) > c) \quad (6)$$

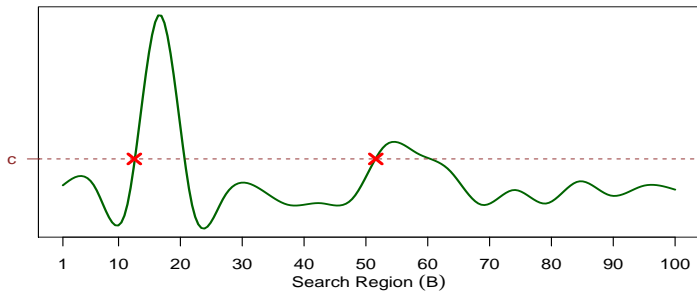
with $c = \max_{\beta_r \in \mathbf{B}_R} LRT(\beta)$.

...which we must calculate/approximate somehow!

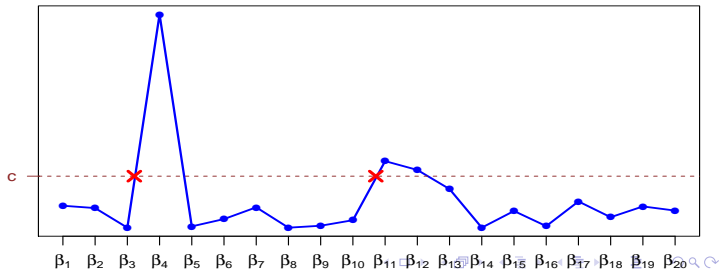
To do so, we first need to introduce the concept of **upcrossings** of the LRT-process $\{LRT(\beta), \beta \in \mathbf{B}\}$.

What do we mean by “upcrossings”?

True LRT-process
under H_0



Discretized version
we deal with
in practice



Approximation of $P(\sup_{\beta \in \mathbf{B}} LRT(\beta) > c)$

- From **Davies, 1987** we have that if $\{LRT(\beta), \beta \in \mathbf{B}\}$ is a “regular” χ_1^2 process, then as $c \rightarrow +\infty$

$$P(\sup LRT(\beta) > c) \approx \frac{P(\chi_1^2 > c)}{2} + \frac{e^{-\frac{c}{2}}}{\sqrt{2\pi}} \int_L^U \kappa(\beta) d\beta \quad \text{Expected \# of upcrossings over } c \text{ of the LRT process under } H_0 \quad (7)$$


- if $c \not\rightarrow +\infty \Rightarrow$ we have an upper bound for $P(\sup LRT(\beta) > c)$.
- $\kappa(\beta)$ is complicated \Rightarrow use the “empirical” version of (7) proposed in

Gross and Vitells, 2010

$$P(\sup LRT(\beta) > c) \approx \frac{P(\chi_1^2 > c)}{2} + \underbrace{e^{-\frac{c-c_0}{2}} E[N(c_0)|H_0]}_{=E[N(c)|H_0]} \quad \text{Expected \# of upcrossings over } c_0 \text{ of the LRT process under } H_0 \quad (8)$$

- where $c_0 \ll c$ and $E[N(c_0)|H_0]$ is estimated using (few) Bootstrap simulations.

For more details and an alternative approach to the problem, check out:

- Algeri S., van Dyk D.A., Conrad J., Brandon, A. *Looking for a Needle in a Haystack? Look Elsewhere! A statistical comparison of approximate global p-values.* Submitted, 2016. 

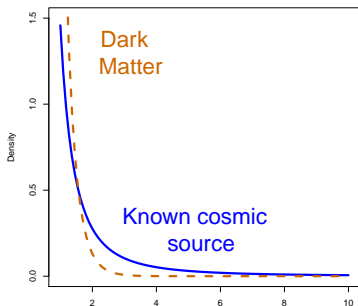
Outline

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Non-nested models comparison in physics

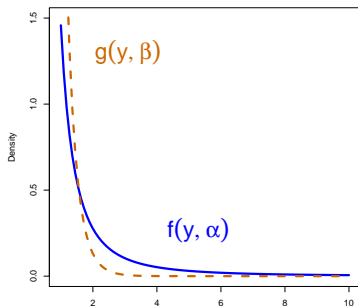
Goal

We would like to distinguish known astrophysics from new signals.



- **E.g.,** Dark Matter.
- We wish to distinguish a dark matter signal from a “fake” signal that mimics it.

The statistical problem



- The model for the known cosmic source is $f(y, \alpha)$;
- The model for the new source is $g(y, \beta)$;
- $f \neq g$ for any α and β .

Is f sufficient to explain the data, or does g provide a better fit?

Problem

f and g are non-nested.

Solutions

Cox, 1961-1962, Atkinson, 1970; etc., Bootstrap, **next two slides..**

Theoretical solutions

Practical solutions

Formulation of the problem

- Consider a comprehensive model which includes $f(y, \alpha)$ and $g(y, \beta)$ as special cases. We have two possibilities:

- Multiplicative form**

$$\propto \{f(y, \alpha)\}^{1-\eta} \{g(y, \beta)\}^{\eta} \quad (9)$$

- Additive form**

$$(1 - \eta)f(y, \alpha) + \eta g(y, \beta) \quad (10)$$

- We prefer (10), it avoids the need to deal with the normalizing constant.

Thus, considering the model in (10) we test

$$H_0 : \eta = 0 \quad \text{versus} \quad H_1 : \eta > 0$$

- To exclude intermediate values of η we can interchange the roles of the hypotheses and test

$$H_0 : \eta = 1 \quad \text{versus} \quad H_1 : \eta < 1.$$

From a new formulation to a well known problem

Model:

$$(1 - \underbrace{\eta}_{\substack{\text{Tested} \\ \text{on the} \\ \text{boundary}}}) f(y, \alpha) + \eta \underbrace{g(y, \beta)}_{\substack{\text{Not} \\ \text{defined} \\ \text{under} \\ H_0}} \quad \text{with} \quad 0 \leq \eta \leq 1 \quad (11)$$

Test:

$$H_0 : \eta = 0 \quad \text{versus} \quad H_1 : \eta > 0$$

similar argument for $H_0 : \eta = 1$ versus $H_1 : \eta < 1$

Note!

These are precisely the same issues we encounter when detecting new particles, i.e., when testing one hypothesis multiple times

⇒ **we already have a solution!**

For more details, check out:

- Algeri S., Conrad J., van Dyk D.A. *A method for comparing non-nested models with application to astrophysical searches for new physics*. MNRAS: Letters, 2016.
- Algeri S., R package 'NONnest', 2015.

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Multiple hypothesis testing - Framework

Also in this case:

- We define a grid \mathbf{B}_R of R β_r values over the energy spectrum \mathbf{B} .
- $\forall \beta_r \in \mathbf{B}_R$ calculate $LRT(\beta_r)$.

However, now we have:

Many ~~sub~~-alternatives...

We have many alternative hypothesis

$H_{11}, \dots, H_{1r}, \dots, H_{1R}$, one for each value $\beta_r \in \mathbf{B}_R$.

...many test statistics...

$\forall \beta_r \in \mathbf{B}_R$ we have one test statistics $LRT(\beta_r)$, and such that $LRT(\beta_r) \sim \frac{1}{2}\chi_1^2 + \frac{1}{2}\delta(0)$ asymptotically.

...many p-values!

$\forall \beta_r \in \mathbf{B}_R$ we have $p_r = \frac{P(\chi_1^2 > LRT(\beta_r))}{2}$.

Local p-values and type I error

- We have an ensemble of R local p-values $p_1, \dots, p_r, \dots, p_R$.
- The smallest, names p_L is then compared with the target probability of type I error α_L .
- But what is α_L if we want to claim a discovery at 5σ ?

Global and local probability of false detection

α_L = specific probability of false detection for each of the R

\neq

α_G = probability of having at least one false detection over the whole ensemble of R tests.

\Rightarrow we must correct p_L accordingly

Local p-values corrections

- **If the R tests were independent**

$$\alpha_G = 1 - (1 - \alpha_L)^R \quad \Rightarrow \quad p_G = 1 - (1 - p_L)^R \quad (12)$$

E.g.: Suppose we are conducting $R = 50$ simultaneous test, each of them at 5σ

$$\alpha_L = 1 - \Phi(5) \quad \Rightarrow \quad \text{by (11):} \quad \alpha_G = 1 - \Phi(4.18)$$

i.e., $\frac{\alpha_G}{\alpha_L} \approx 50$.

- **If the R tests were dependent** (which is generally the case)

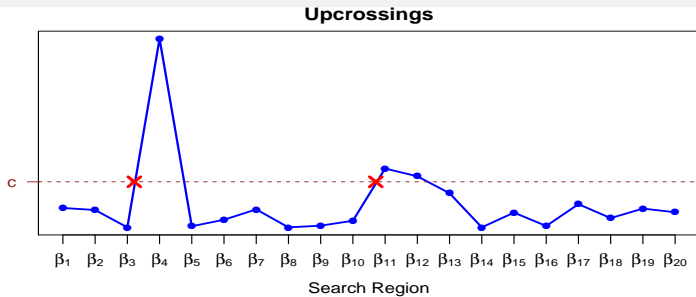
$$\alpha_G \leq R\alpha_L \quad \Rightarrow \quad p_{BF} = Rp_L \quad \text{Bonferroni's correction} \quad (13)$$

Outline

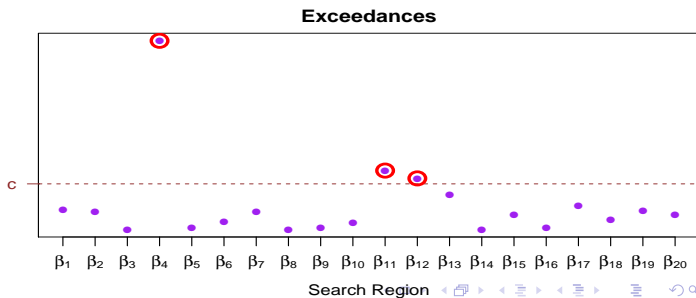
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Upcrossings and Exceedances

Discretized
LRT-process
under H_0



Multiple LRTs
under H_0



Why are we interested in the Exceedances?

We can identify situations where the average number of exceedances under H_0 , namely $E[N_c^*|H_0]$, and the average number of upcrossings under H_0 , $E[N_c|H_0]$ are approximately equal.

- We will soon see two conditions we need for this to happen.
- For now let's focus on $E[N_c^*|H_0]$:

$$E[N_c^*|H_0] = \sum_{r=1}^R 1 \cdot P(LRT(\beta_r) > c)$$

under H_0 , $\forall \beta_r \in \mathbf{B}_R$,

$$LRT(\beta_r) \sim \frac{1}{2}\chi_1^2 + \frac{1}{2}\delta(0) \text{ asymptotically}$$

$$= \sum_{r=1}^R \frac{P(\chi_1^2 > c)}{2} = R \frac{P(\chi_1^2 > c)}{2} = R p_L = p_{BF} \quad \text{Bonferroni's correction!}$$

Two sides of the same coin

What the previous slide is telling us is that, if

$$\underbrace{E[N_c|H_0]}_{\substack{\text{Expected} \\ \# \text{ upcrossings} \\ \text{under } H_0}} \approx \underbrace{E[N_c^*|H_0]}_{\substack{\text{Expected} \\ \# \text{ exceedances} \\ \text{under } H_0}} = \underbrace{p_{BF}}_{\substack{\text{Bonferroni's} \\ \text{correction}}} \quad (14)$$

and $\exists \lambda$ s.t as $c \rightarrow +\infty$ ($p_L \rightarrow 0$) and $R \rightarrow +\infty$

$$E[N_c|H_0] \approx E[N_c^*|H_0] = p_{BF} \rightarrow \lambda$$

then, for R and c large we have

$$\underbrace{P(\sup LRT(\beta) > c)}_{\text{Global p-value}} \approx \underbrace{\frac{P(\chi_1^2 > c)}{2}}_{\substack{\rightarrow 0 \\ \text{as } c \rightarrow +\infty}} + E[N_c|H_0] \quad (15)$$

$$\approx E[N_c|H_0] \approx E[N_c^*|H_0]$$

$$\approx p_{BF} \quad \text{Bonferroni adjusted local p-value}$$

This means that if $E[N_c|H_0] \approx E[N_c^*|H_0]$, then testing one hypothesis multiple times and multiple hypothesis testing will lead to approximately the same inference. (But, since the latter is much quicker than the former, I might gain in computing time.)

When do we have $E[N_c|H_0] \approx E[N_c^*|H_0]$?

To guarantee $E[N_c|H_0] \approx E[N_c^*|H_0]$ (as $c \rightarrow +\infty$), we need the following two conditions to be satisfied:

1 Long range independence

$$|F_{1,\dots,r,r+1,\dots,r+k} - F_{1,\dots,r}F_{r+1,\dots,r+k}| \leq q(r) \quad (16)$$

where $F(\cdot)$ is the cdf of $LRT(\beta_r)$, $\forall \beta_r \in \mathbf{B}_R$, and $q(r)$ is a function such that $q(r) \rightarrow 0$ as $r \rightarrow \infty$.

This condition implies that independence is achieved for distant points β_r of the (discretized) energy/mass spectrum.

2 Local dependence

$$\limsup R \sum_{r=2}^{[R/l]} P(LRT(\beta_1) > c, LRT(\beta_r) > c) \rightarrow 0 \quad \text{as } l \rightarrow +\infty \quad (17)$$

where $F(\cdot)$ be the cdf of $LRT(\beta_r)$, $\forall \beta_r \in \mathbf{B}_R$,

This condition excludes the chance of clustering of the upcrossings of the LRT-process.

How to assess if these two conditions hold?

- Let the model of reference be $(1 - \eta)f(y, \alpha) + \eta g(y, \beta_r)$, and let $l(\eta|\alpha, \beta_r, y)$ be its log-likelihood.
- $\forall \beta_r$ the score function evaluated at H_0 is $S(\beta_r) = \frac{\partial l(\eta|\alpha, \beta_r, y)}{\partial \eta} \Big|_{\eta=0}$
 \Rightarrow the score process under H_0 is $\{S(\beta_r), \beta_r \in \mathbf{B}_r\}$
- with covariance function is $\text{cov}(S(\beta_r), S(\beta_t)) = \int \frac{g(y, \beta_r)g(y, \beta_t)}{f(y, \alpha)} \partial y - 1$

$$S^*(\beta_r) = \frac{S(\beta_r)}{\sqrt{\text{cov}(S(\beta_r), S(\beta_r))}} \quad (18)$$

A sufficient condition on $S^*(\beta_r)$ (Berman's condition)

If the covariance function of $S^*(\beta_r)$ satisfies

$$\sup_{|\beta_r - \beta_t| > \tau} |\text{cov}(S^*(\beta_r), S^*(\beta_t))| \log(\tau) \rightarrow 0 \quad \text{as } \tau \rightarrow +\infty \quad (19)$$

then **long range independence** and **local independence** hold on both the normalized score and the LRT processes.

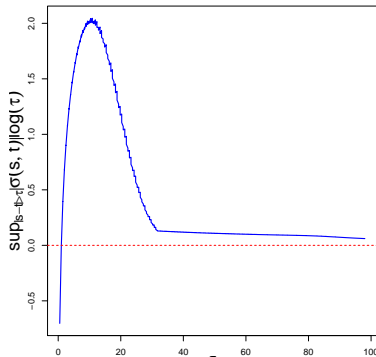
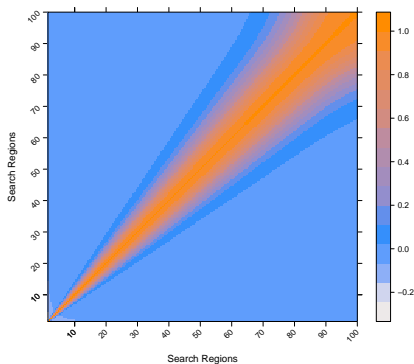
Example

Consider a power-law distributed background with index ψ and a Gaussian signal with dispersion proportional to the signal location.

The full model is

$$(1 - \eta) \frac{1}{k_\psi y^{\psi+1}} + \frac{\eta}{k_{M_\chi}} \exp\left\{-\frac{(y - M_\chi)^2}{0.02 M_\chi^2}\right\} \quad (20)$$

with k_ψ and k_{M_χ} normalizing constants, $y \in [1; 100]$, $\psi = 1.4$ and $M_\chi \in [1; 100]$.



Realistic data analysis

We simulated observation of monochromatic feature by the Fermi Large Area Telescope (LAT).

- 2391 events from an astrophysical background corresponding to isotropic emission following a spectral power-law with index 2.4, i.e., $\psi = 1.4$.
- 64 events from a Gaussian signal with mass of 35 GeV.
- 80 energy bins, spaced equally from 10-350 GeV.

Method	Signal Location	Signal Strength	Sig.
Unadjusted local	35.82	0.042	5.920 σ
Bonferroni adj. local	35.82	0.042	5.152 σ
Gross & Vitells	35.82	0.042	5.192 σ

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What can we do more?

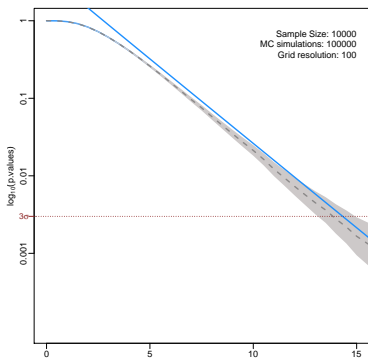
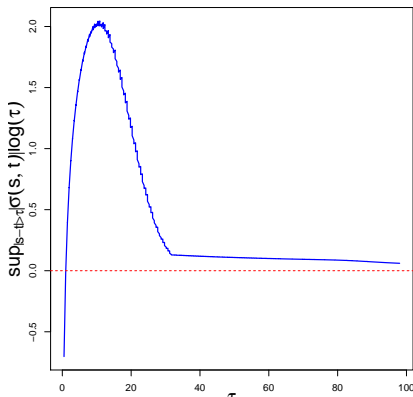
- Berman's condition is not only a sufficient condition to guarantee asymptotic equivalence between testing one hypothesis multiple times and multiple hypothesis testing.
- Indeed, it can be used as diagnostic tool to assess the validity of the Davies (1987) and Gross and Vitells (2010) approximations for the global p-value $P(\sup LRT(\beta) > c)$.
- Several cases can be identified and additional conditions, in addition to long range independence and local independence, are needed.
- But we still have to refine the details...
- ...however, we already can take a look at some examples.

A case where everything works nicely

Considering again the Power Law background + Gaussian signal example:

$$(1 - \eta) \frac{1}{k_\psi y^{\psi+1}} + \frac{\eta}{k_{M_\chi}} \exp\left\{-\frac{(y - M_\chi)^2}{0.02 M_\chi^2}\right\} \quad (21)$$

with k_ψ and k_{M_χ} normalizing constants, $y \in [1; 100]$, $\psi = 1.4$ and $M_\chi \in [1; 100]$.

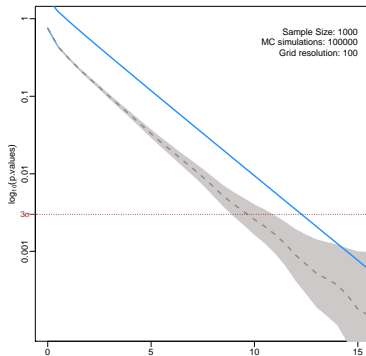
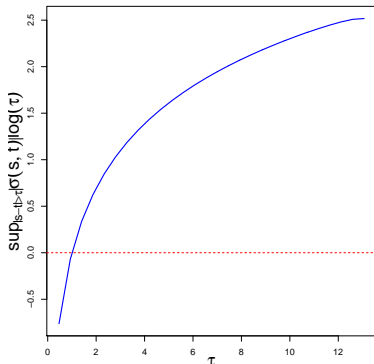


A non-ideal case

Suppose we want to distinguish between Pulsar Spectrum and Dark Matter. The full model is:

$$(1 - \eta) \frac{\exp\{-y^2\}}{k_\rho y^\rho} + \frac{\eta \exp\{-7.8 \frac{y}{\phi}\}}{k_\phi y^{1.5}} \quad (22)$$

with k_ρ and k_ϕ normalizing constants, $y \in [1; 15]$, $\rho = 4/3$ and $\phi \in [1; 15]$.

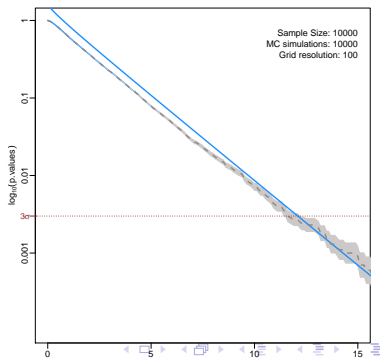
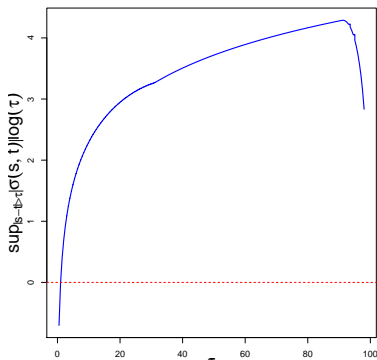


A case somewhere in between

Suppose we want to distinguish between a Power Law distributed cosmic source and Dark Matter. The full model is:

$$(1 - \eta) \frac{1}{k_\psi y^{\psi+1}} + \frac{\eta \exp\{-7.8 \frac{y}{\phi}\}}{k_\phi y^{1.5}} \quad (23)$$

with k_ψ and k_ϕ normalizing constants, $y \in [1; 100]$, $\psi = 1.4$ and $\phi \in [1; 100]$.



...in the “next episode” ...

Work in progress and (immediate) future goals:

- We would like to provide a formal explanation of cases where the global p-value approximations do and do not work.
- We would like to provide precise indications on how to spot these cases.
- We would like to exploit the information on the dependence structure of the underlying processes to improve, if possible, the global p-value approximations discussed in this talk.

All this will be discussed in:

Algeri S., van Dyk D.A., Conrad J. *Testing one hypothesis multiple times*. In preparation, 2016.

(Hopefully, available on ArXiv by the end of the summer.)

Thank you for listening!

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- Algeri S., van Dyk D.A., Conrad J., Brandon, A. *Looking for a Needle in a Haystack? Look Elsewhere! A statistical comparison of approximate global p-values*. Submitted, 2016.
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